# The volume of pseudo-effective line bundles and partial Okounkov bodies 

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## Table of Contents

## (1) Background

## (2) The volume of Hermitian pseudo-effective line bundles

(3) Partial Okounkov bodies

4 Applications

## Singular Hermitian metrics on line bundles

- $X$ : complex projective manifold of dimension $n$.
- $L$ : a holomorphic line bundle on $X$.


## Singular Hermitian metrics on line bundles

- $X$ : complex projective manifold of dimension $n$.
- $L$ : a holomorphic line bundle on $X$.


## Singular Hermitian metric on $L$

Locally on $U \subseteq X$, we can identify $L$ with the trivial bundle $U \times \mathbb{C}$, a singular Hermitian is given by $e^{-\varphi}$, where $\varphi$ is a locally integrable function on $U$ : given any local section $s$ of $L$ over $U$, the square norm of $s$ at $x \in U$ is given by

$$
|s|_{\phi}^{2}(x)=|s(x)|^{2} e^{-\varphi(x)}
$$

This definition can be globalized in the obvious manner, which gives the notion of a (singular) Hermitian metric $\phi$ on $L$.

## Curvature form

Let $\mathrm{dd}^{\mathrm{c}} \phi$ denote the curvature form/current of the metric $\phi$ : locally represent the metric $\phi$ by weights $e^{-\varphi}$, then

$$
\mathrm{dd}^{\mathrm{c}} \phi:=\frac{\mathrm{i}}{2 \pi} \partial \bar{\partial} \varphi
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as currents.

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as currents.

## Proposition

$\operatorname{dd}^{\mathrm{c}} \phi$ is in the first Chern class $c_{1}(L)$ of $L$.
This explains our normalization of $\mathrm{dd}^{\mathrm{c}}$.

## Big line bundles

## Definition

We say $L$ is big if there exists $c>0$ so that

$$
h^{0}\left(X, L^{k}\right)>c k^{n}
$$

for any $k \gg 1$.

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$$

for any $k \gg 1$.

## Example

All ample line bundles are big. In fact, by Riemann-Roch theorem

$$
h^{0}\left(X, L^{k}\right)=\frac{\left(L^{n}\right)}{n!} k^{n}+\mathcal{O}\left(k^{n-1}\right)
$$

## Pseudo-effective line bundles

Fix a Kähler form $\omega$ on $X$.

## Theorem (Demailly)

The line bundle $L$ is big if and only if $L$ admits a singular upper semi-continuous Hermitian metric $\phi$ so that $\mathrm{dd}^{\mathrm{c}} \phi$ is a Kähler current, namely,

$$
\operatorname{dd}^{\mathrm{c}} \phi \geq \epsilon \omega
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for some $\epsilon>0$.

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## Definition

$L$ is said to be pseudo-effective (psef) if there is an upper semi-continuous singular Hermitian metric $\phi$ on $L$ with $\mathrm{dd}^{\mathrm{c}} \phi \geq 0$ (a positive/psh metric). A pair $(L, \phi)$ consisting of a psef line bundle $L$ and a positive metric $\phi$ on $L$ is called a Hermitian pseudo-effective line bundle.

## Goal of this talk

There are well-developed theories of psef line bundles: we have (among others)

- An intersection theory (movable intersection).
- A theory of volume.
- A theory of Okounkov bodies.


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There are well-developed theories of psef line bundles: we have (among others)

- An intersection theory (movable intersection).
- A theory of volume.
- A theory of Okounkov bodies.

Each theory has a generalization in the Hermitian psef line bundle setting:

- The theory of non-pluripolar products (Boucksom-Eyssidieux-Guedj-Zeriahi, ...).
- A theory of volume (Darvas-X.).
- The theory of partial Okounkov bodies (X.).

We will explain the last two theories in this talk.

## Table of Contents

## (1) Background

(2) The volume of Hermitian pseudo-effective line bundles

## (3) Partial Okounkov bodies

4 Applications

## The theory of volumes of pseudo-effective line bundles

Let $L$ be a psef line bundle on $X$.

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The volume of $L$ is defined as

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\operatorname{vol}(L):=\lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k}\right)
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The existence of the limit is a non-trivial fact.

## The volume of Hermitian pseudo-effective line bundles

## Example

When $L$ is ample/nef, vol $L=\frac{1}{n!}\left(L^{n}\right)$. More generally, if $L$ is only pseudo-effective, $\operatorname{vol} L=\frac{1}{n!}\left\langle L^{n}\right\rangle$ (movable intersection).

## Proposition

1. The volume of a line bundle depends only on its numerical class: if $c_{1}(L)=c_{1}\left(L^{\prime}\right)$, then $\operatorname{vol} L=\operatorname{vol} L^{\prime}$.
2. The volume is homogeneous: $\operatorname{vol} L^{k}=k^{n} \operatorname{vol} L$.
3. A psef line bundle $L$ is big iff vol $L>0$.

## How to generalize to singular setting?

We want to generalize

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\operatorname{vol} L=\lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k}\right)
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## Definition

The multiplier ideal sheaf of $\phi$ is the (coherent) subsheaf of $\mathcal{O}_{X}$, locally consisting of holomorphic functions $f$ such that $|f|_{\phi}^{2}$ is locally integrable.

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We can thus define

$$
\operatorname{vol}(L, \phi)=\lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right)
$$

Existence of the limit is non-trivial!

## Why and how do we study volumes?

## Why?

From the perspective of algebraic geometry, we have the following vanishing theorem:

## Theorem (Nadel-Cao+Darvas-X)

Suppose $\operatorname{vol}(L, \phi)>0$, then

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h^{i}\left(X, K_{X} \otimes L^{k} \otimes \mathcal{J}(k \phi)\right)=0
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for all $i>0, k>0$.
How?
(1) Treat the case of analytic singularities (Bonavero)
(2) Treat the general case by Demailly approximation.

## Bonavero's theorem

Assume that $\phi$ has analytic singularities: locally $\phi$ is of the form $c \log \sum_{i}\left|f_{i}\right|^{2}$ (up to a smooth term), where $f_{1}, \ldots, f_{M}$ are holomorphic functions.

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Theorem (Bonavero)

$$
\lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right)=\frac{1}{n!} \int_{X}\left(\operatorname{dd}^{\mathrm{c}} \phi\right)^{n}
$$

Here $\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}$ is the non-pluripolar product.

## Non-pluripolar products

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When $\phi$ is smooth, $\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}$ is the usual wedge product of forms. When $\phi$ is bounded, $\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}$ coincides with the Bedford-Taylor product. For general $\phi,\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}$ is a measure on $X$ that puts no mass on pluripolar sets.

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The theory of non-pluripolar products is not the only extension of Bedford-Taylor theory (Dinh-Sibony products, residue currents, etc.). But non-pluripolar products have two advantages: $\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}$ can be defined for all $\phi$ and

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## Theorem (Witt Nyström)

If $\varphi \leq \psi$, then

$$
\int_{X}\left(\operatorname{dd}^{\mathrm{c}} \varphi\right)^{n} \leq \int_{X}\left(\mathrm{dd}^{\mathrm{c}} \psi\right)^{n}
$$

## Quasi-equisingular approximations

Assume that $\operatorname{dd}^{\mathrm{c}} \phi$ is a Kähler current. We can find a decreasing sequence ( $\phi^{j}$ ) of psh metrics on $L$ of analytic singularities converging to $\phi$ such that for any $\delta>0, k>0$, there is $j_{0}>0$ so that for $j \geq j_{0}$,

$$
\mathcal{J}\left((1+\delta) k \phi^{j}\right) \subseteq \mathcal{J}(k \phi) \subseteq \mathcal{J}\left(k \phi^{j}\right) .
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## To remember

Quasi-equisingular approximations are optimal analytic approximations of psh metrics in the sense of multiplier ideal sheaves.

## From Bonavero to general potentials

By approximation,

$$
\begin{aligned}
\varlimsup_{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right) & \leq \lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}\left(k \phi^{j}\right)\right) & & \text { as } \phi \leq \phi^{j} \\
& =\frac{1}{n!} \int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi^{j}\right)^{n} & & \text { Bonavero } .
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By Witt Nyström, $\int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi^{j}\right)^{n}$ is decreasing in $j$.

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$$

In general,

$$
\lim _{j \rightarrow \infty} \int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi^{j}\right)^{n} \neq \int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}
$$

What is the limit? To answer this question, we need to introduce $\mathcal{J}$-model potentials.

## I-model potentials

Given a smooth form $\theta \in c_{1}(L)$, then $\phi$ can be identified with $\varphi \in \operatorname{PSH}(X, \theta)$.

## Definition

Define

$$
P[\varphi]_{\mathcal{J}}:=\sup \{\psi \leq 0: \mathcal{J}(k \varphi)=\mathcal{J}(k \psi)\} .
$$

Equivalently,

$$
P[\varphi]_{\mathcal{J}}:=\sup \left\{\psi \leq 0: \nu\left(\pi^{*} \varphi, y\right)=\nu\left(\pi^{*} \psi, y\right)\right\}
$$

for any birational model $\pi: Y \rightarrow X$ and any $y \in Y$.
A potential $\varphi$ with $\varphi=P[\varphi]_{\mathcal{J}}$ is called an $\mathcal{J}$-model potential.

## I-model potential

Theorem (Darvas-X)

$$
\lim _{j \rightarrow \infty} \int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi^{j}\right)^{n}=\int_{X}\left(\theta+\mathrm{dd}^{\mathrm{c}} P[\varphi]_{\mathcal{J}}\right)^{n}
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## I-model potential

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As a consequence,

$$
\varlimsup_{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right) \leq \frac{1}{n!} \int_{X}\left(\theta+\operatorname{dd}^{\mathrm{c}} P[\varphi]_{\mathcal{J}}\right)^{n}
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## I-model potential

## Theorem (Darvas-X)

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$$

The reverse inequality is much harder. But the main idea is the same.

## Theorem

When $\operatorname{dd}^{\mathrm{c}} \phi$ is a Kähler current, we have

$$
\operatorname{vol}(L, \phi)=\lim _{k \rightarrow \infty} \frac{1}{k^{n}} h^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right)=\frac{1}{n!} \int_{X}\left(\theta+\operatorname{dd}^{\mathrm{c}} P[\varphi]_{\mathcal{J}}\right)^{n}
$$

## Main theorem

When $\operatorname{dd}^{\mathrm{c}} \phi$ is not a Kähler current, a trick allows us to reduce to the special case: there is always a potential $\psi \leq \phi$ such that $\operatorname{dd}^{\mathrm{c}} \psi$ is a Kähler current. We use $(1-\epsilon) \phi+\epsilon \psi$ to approximate $\phi$.

## Theorem

Let $(L, \phi)$ be a Hermitian psef line bundle on $X$, under the identification $\phi$ with $\varphi \in \operatorname{PSH}(X, \theta), \operatorname{vol}(L, \phi)$ is well-defined and

$$
\operatorname{vol}(L, \phi)=\frac{1}{n!} \int_{X}\left(\theta+\operatorname{dd}^{\mathrm{c}} P[\varphi]_{\mathcal{J}}\right)^{n}
$$

## Proof recap

The main steps in the proof can be sketched as follow:

- Treat the case of analytic singularities.
- Treat the case of Kähler currents by approximation.
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## Proof recap

The main steps in the proof can be sketched as follow:

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This general reduction scheme will be applied in other contexts as well.

## Table of Contents

## (1) Background

## (2) The volume of Hermitian pseudo-effective line bundles

(3) Partial Okounkov bodies

4 Applications

## Okounkov bodies

Fix a big line bundle $L$ on $X$.
Okounkov body is a machinery that allows one to study numerical properties of $L$ using convex geometry.

## Definition

An admissible flag is a sequence

$$
X=Y_{0} \supseteq Y_{1} \supseteq \cdots \supseteq Y_{n}
$$

of subvarieties of $X, Y_{i}$ of codimension $i$ such that $Y_{i}$ is smooth at the point $Y_{n}$.

## Example

When $X$ is a Riemann surface, an admissible flag is $X \supseteq\{$ point $\}$.

## Okounkov bodies

Fix an admissible flag $\left(Y_{\bullet}\right)$. Given any $s \in \mathbb{C}(X)^{\times}$, we can define the valuation of $s$ along $\left(Y_{\bullet}\right)$ :

$$
\nu(s)=\left(\nu_{1}(s), \ldots, \nu_{n}(s)\right) \in \mathbb{Z}^{n}
$$

Roughly speaking, $\nu_{1}(s)$ is the order of vanishing of $s$ along $Y_{1}$. If $\nu_{1}(s)=0, \nu_{2}(s)$ is the order of vanishing of $\left.s\right|_{Y_{1}}$, etc..

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Set

$$
\Gamma(L):=\left\{(\nu(s), k) \in \mathbb{Z}^{n+1}: s \in H^{0}\left(X, L^{k}\right)\right\}
$$

Then $\Gamma(L)$ is a semi-group in $\mathbb{Z}^{n+1}$ and generates a closed convex cone $C$ in $\mathbb{R}^{n+1}$. We define the Okounkov body of $L$ as

$$
\Delta(L):=C \cap\left\{(x, 1): x \in \mathbb{R}^{n}\right\}
$$

## Example

To understand the abstract definition, take $X=\mathbb{P}^{1}$. The flag is $X \supseteq\{0\}$. Given any rational function $f, \nu(f)$ is the order of vanishing of $f$ at 0 .

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$$
\Gamma(L)=\left\{(a, b) \in \mathbb{Z}^{2}: 0 \leq a \leq b\right\}
$$

## Example



## How to understand Okounkov bodies?

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When $n=1$, flag is just a point and $\Delta(L)=[0, \operatorname{deg} L]$.

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## Convex body

A convex body is a non-empty compact convex subset of $\mathbb{R}^{n}$.
In general, $\Delta(L)$ is a convex body, but not a polytope!

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## Convex body

A convex body is a non-empty compact convex subset of $\mathbb{R}^{n}$.
In general, $\Delta(L)$ is a convex body, but not a polytope! Okounkov bodies is a way to apply techniques in toric geometry to non-toric varieties.

## Why are Okounkov bodies important?

Theorem (Lazarsfeld-Mustață, Kaveh-Khovanskii)

$$
\frac{1}{n!} \operatorname{vol}(L)=\operatorname{vol} \Delta(L)
$$

## Why are Okounkov bodies important?

## Theorem (Lazarsfeld-Mustață, Kaveh-Khovanskii) <br> $$
\frac{1}{n!} \operatorname{vol}(L)=\operatorname{vol} \Delta(L)
$$

## Theorem (Lazarsfeld-Mustață, Jow)

$\Delta(L)$ depends only on the numerical class of $L$. Moreover, the information of $\Delta(L)$ for all admissible flags determines the numerical class of $L$ uniquely.

In particular, we can (in principle) deduce all numerical information of $L$ from $\Delta(L)$.

## Partial Okounkov bodies

Based on the same ideas, we want to define Okounkov bodies $\Delta(L, \phi)$ of Hermitian psef line bundles $(L, \phi)$ with similar properties.

## Partial Okounkov bodies

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## Partial Okounkov bodies

Based on the same ideas, we want to define Okounkov bodies $\Delta(L, \phi)$ of Hermitian psef line bundles $(L, \phi)$ with similar properties.
In contrast to concluding numerical information of $L$, we want to get information about the singularities of $\phi$.

## Theorem (X.)

There is a canonical convex body $\Delta(L, \phi) \subseteq \Delta(L)$ associated to each Hermitian psef line bundle with $\int_{X}\left(\mathrm{dd}^{\mathrm{c}} \phi\right)^{n}>0$. The assignment is continuous in $\phi$.
Two psh metrics (of positive masses) $\phi$ and $\phi^{\prime}$ have the same partial Okounkov bodies iff $\phi$ and $\phi^{\prime}$ are $\mathcal{J}$-equivalent: $\mathcal{J}(k \phi)=\mathcal{J}\left(k \phi^{\prime}\right)$ for all $k$.

## How to construct partial Okounkov bodies?

Recall our general construction scheme:

- Treat the case of analytic singularities.
- Treat the case of Kähler currents by approximation.
- Reduce the general case to the case of Kähler currents.


## Partial Okounkov bodies-Analytic singularities

When $\phi$ has analytic singularities.
Assume for simplicity that $\phi$ has logarithmic singularities along some normal crossing divisor $D: \phi$ is locally $\log |s|^{2}$, where $s$ is the defining equation of $D$.

## Partial Okounkov bodies-Analytic singularities

When $\phi$ has analytic singularities.
Assume for simplicity that $\phi$ has logarithmic singularities along some normal crossing divisor $D: \phi$ is locally $\log |s|^{2}$, where $s$ is the defining equation of $D$. In this case,

$$
H^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right)=H^{0}\left(X, L^{k} \otimes \mathcal{O}(-k D)\right)
$$

So

$$
\bigoplus_{k} H^{0}\left(X, L^{k} \otimes \mathcal{J}(k \phi)\right)=\bigoplus_{k} H^{0}\left(X, L^{k} \otimes \mathcal{O}(-k D)\right)
$$

## Partial Okounkov bodies-Analytic singularities

Define

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\Delta(L, \phi):=\Delta(L-D)+a,
$$

where $a=\nu(D) \in \mathbb{R}^{n}$ is a fixed vector.

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where $a=\nu(D) \in \mathbb{R}^{n}$ is a fixed vector.
The general case of analytic singularities can be handled in a similar fashion. This finishes Step 1 in our scheme.

## Partial Okounkov bodies-Kähler currents

Step 2 in our scheme is the quasi-equisingular approximation $\left(\phi^{j}\right)$ of $\phi$. We want to define

$$
\Delta(L, \phi):=\lim _{j \rightarrow \infty} \Delta\left(L, \phi^{j}\right)
$$

But what is the topology on convex bodies $\Delta(\bullet)$ ?

## Partial Okounkov bodies-Kähler currents

Step 2 in our scheme is the quasi-equisingular approximation $\left(\phi^{j}\right)$ of $\phi$.
We want to define

$$
\Delta(L, \phi):=\lim _{j \rightarrow \infty} \Delta\left(L, \phi^{j}\right)
$$

But what is the topology on convex bodies $\Delta(\bullet)$ ?
The solution is given by the Hausdorff topology.

## Definition

Given two convex bodies $K_{1}, K_{2} \subseteq \mathbb{R}^{n}$, define their Hausdorff metric as

$$
d_{n}\left(K_{1}, K_{2}\right):=\max \left\{\sup _{x_{1} \in K_{1}} \inf _{x_{2} \in K_{2}}\left|x_{1}-x_{2}\right|, \sup _{x_{2} \in K_{2}} \inf _{x_{1} \in K_{1}}\left|x_{1}-x_{2}\right|\right\}
$$

## Partial Okounkov bodies-General case

Step 3 in our scheme is carried our in a similar way as Step 2. We can reduce general potentials to those whose curvature currents are Kähler currents easily.

## Partial Okounkov bodies-General case

Step 3 in our scheme is carried our in a similar way as Step 2. We can reduce general potentials to those whose curvature currents are Kähler currents easily.
The construction seems to depend on a lot of choices. But we can prove

## Theorem

The convex body $\Delta(L, \phi)$ is independent of all choices we made! It depends on $(L, \phi)$ only through the current $\mathrm{dd}^{\mathrm{c}} \phi$.

In other words, $\Delta(L, \phi)$ is a numerical invariant of $(L, \phi)$.

## Why do we care about partial Okounkov bodies?

As we mentioned,

## Theorem

Let $\phi, \phi^{\prime}$ be two psh metrics on $L$ with positive masses. Then
$\Delta(L, \phi)=\Delta\left(L, \phi^{\prime}\right)$ for all admissible flags (on all birational models of $X$ ) iff $\phi \sim_{\mathcal{J}} \phi^{\prime}$.

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In other words,

## Principle

Partial Okounkov bodies are universal invariants of the singularities of a psh metric.

## Dictionary

We have a the following dictionary between psef line bundles and Hermitian psef line bundles.

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- Many others


## How do we study/classify general psh singularities?

We discuss a general classification of the singularities of psh metrics of positive volumes.

- Classification according to $\mathcal{J}$-singularity types. Each class admits an $\mathcal{J}$-model representative, which can be studied using partial Okounkov bodies, namely by convex geometry.


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- In each $\mathcal{J}$-class, we can subdivide according to the model types, namely, according to the behaviors of their Monge-Ampère type masses. In each model class, the Monge-Ampère equation is uniquely solvable, hence we can study each class using Monge-Ampère measures.


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- In each model class, we can subdivide according to singularities types.


## Table of Contents

## (1) Background

## (2) The volume of Hermitian pseudo-effective line bundles

(3) Partial Okounkov bodies

4 Applications

## A few applications

In pluripotential theory, our result yields a very general Bergman kernel convergence result: given a singularity type [ $\varphi$ ], a continuous metric $\psi$ on $L$, it is possible to define partial Bergman measures $\beta_{[\varphi]}^{k}(\psi)$, generalizing the usual Bergman measure construction.

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We proved that these partial Bergman measures converge weakly and the limit can be described explicitly. This result can be regarded as a local version of our volume theorem.

## A few applications

When $L$ is ample, our theory allows us to understand the non-Archimedean space $\mathcal{E}^{1}\left(L^{\text {an }}\right)$ defined by Boucksom-Jonsson in terms of the space of geodesic rays in the Archimedean (usual) setting. In particular, we have a complete description of maximal geodesic rays. One can also construct Duistermaat-Heckman measures of elements in $\mathcal{E}^{1}\left(L^{\mathrm{an}}\right)$, which opens the gate to the study of non-Archimedean space $\mathcal{E}^{p}\left(L^{\mathrm{an}}\right)$.

## A few applications

Very recently, Botero-Burgos Gil-Holmes-De Jong made use of our results to study the canonical metric on the line bundle of Siegel-Jacobi forms: Let $A_{g, N}$ be the fine moduli space of PPAV of dimension $g$ and level $N$. Consider the universal Abelian variety $U_{g, N}$ on $A_{g, N}$ and the line bundle $\mathcal{L}_{k, m}$ of Siegel-Jacobi forms on $U_{g, N}$ of weight $k$ and index $m$. The latter is endowed with a natural metric $h$. They applied our results to study the behaviour of $h$ on the compactification of $U_{g, N}$.

## Thank you for your attention!

